

# Lecture 8: Inviscid Burgers equation

①

Equation arises as geodesic on  $\text{Diff}(M)$ :

$$A(x) = \frac{1}{2} \int_{t_1}^{t_2} \int_M |\dot{X}_t(a)|^2 da dt$$

$$\delta A(x) = 0 \iff \ddot{X}_t = 0$$

The equation for  $M \subseteq \mathbb{R}^n$  (without boundary) is:

$$\text{Initial value problem} \begin{cases} \ddot{X}_t(a) = 0 \\ \dot{X}_0(a) = u_0(a) \\ X_0(a) = a \in M \end{cases}$$

is called the Burgers equation. If  $M$  is flat, then trajectories  $X_t$  themselves sweep out geodesics on  $M$ . Provided  $X$  is a diffeomorphism;

Eulerian form:

$$\dot{X}_t(a) = u(X_t(a), t)$$

$$u: M \times \mathbb{R} \rightarrow \mathbb{R}^d$$

$$\begin{aligned} \partial_t u + u \cdot \nabla u &= 0 \\ u|_{t=0} &= u_0 \end{aligned}$$

# Two point problem for Burgers (remark) (2)

Fix  $\phi \in \text{Diff}(M)$ . Burgers geodesic between  $\phi$  and  $\text{id}$ :

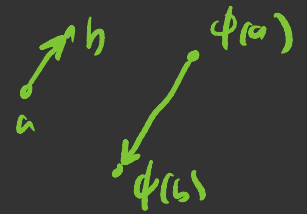
Two-point problem  $\begin{cases} \ddot{X}_t(a) = 0 \\ X_0(a) = a \\ X_1(a) = \phi(a) \end{cases} \iff X_t(a) = (1-t)a + t\phi(a)$

Suppose  $\exists t_* > 0$  s.t.  $X_{t_*}(a) = X_{t_*}(b)$  for  $a \neq b$ . This happens iff  $a - b = \frac{-t_*}{1-t_*} (\phi(a) - \phi(b))$ .

Note  $f(t) = \frac{-t}{1-t} \in (-\infty, 0)$

Remark: in 1d, say  $a > b$ , then  $\phi(a) > \phi(b)$

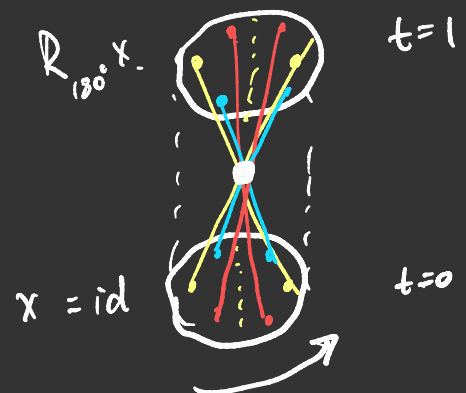
requires anti-alignment



Thus  $\frac{a-b}{\phi(a)-\phi(b)} > 0$  and so there cannot be collisions.

In 1d, the two point problem is always solvable

Remark in 2D: solid body has collisions once a half turn is made.



# Singularity formation v.s. Global wellposedness

3

$$\dot{X}_t(a) = 0 \Rightarrow X_t(a) = a + t u_0(a)$$

Introducing  $X_t(a) = u(X_t(a), t)$ , we see

$$\frac{d}{dt} u(X_t(a), t) = 0 \Rightarrow u(X_t(a), t) = u_0(a)$$

Combining, we have

$$u(a + t u_0(a), t) = u_0(a)$$

Differentiating in the label "a",

$$\nabla_a [u(a + t u_0(a), t)] = \nabla_a u_0(a)$$

Thus

$$(\mathbf{I} + t \nabla_a u_0(a)) \cdot (\nabla_x u)(a + t u_0(a), t) = \nabla_a u_0(a)$$

which implies

$$\nabla_x u(a + t u_0(a), t) = (\mathbf{I} + t \nabla_a u_0(a))^{-1} \nabla_a u_0(a)$$

$$\nabla_x u(a + t u_0(a), t) = (\mathbb{I} + t \nabla_a u_0(a))^{-1} \nabla_a u_0(a)$$

Remarks: Let  $A_0 = \nabla_a u_0(a)$ ,  $A(t) = \nabla_x u(a + t u_0(a), t)$ .

① If  $A_0$  has no real eigenvalues (excepting possibly 0) then the solution to Burgers exists for all time  $t \in \mathbb{R}$ .

Impossible if  $d=1$ !

②  $A(t)$  is determined for all time  $t \geq 0$  (and  $t < 0$ ) if  $A_0$  has no negative (corresp. no positive) eigenvalues.

Compressing in one or other time direction.

③ Incompressible motions!

Such are called asymptotic geodesics, since they are geodesic both on the constrained space  $D_\mu(M)$  as well as  $P(M)$

↑ submanifold of  $D(M)$ .

# Asymptotic geodesics

(5)

Recall  $X_t(a) = a + t u_0(a)$

$$\nabla X_t(a) = I + t A_0 \quad (A_0 = \nabla_a u_0(a))$$

If  $X_t \in D_\mu(m)$ , then  $\det \nabla X_t(a) = 1$  for all  $t \in \mathbb{R}$

In this case, we must have

$$\det(I + t A_0) = 1 \quad \text{for all } t \in \mathbb{R}$$

Lemma: let  $N \in M^{n,n}(\mathbb{R})$ . Then

$$\det(I + tN) = 1 \quad \forall t \in \mathbb{R} \iff N \text{ nilpotent}$$

Proof: Assuming  $\det(I + tN) = 1$  we have:

$$0 = \frac{d}{dt} \det(I + tN) = \text{tr}((I + tN)^{-1} N) \quad \forall t.$$

Evaluating at  $t=0$ , we find  $\text{tr} N = 0$ . Differentiating again:

$$0 = \frac{d}{dt} \text{tr}((I + tN)^{-1} N) \Big|_{t=0} = \text{tr}((I - tN)^{-1} N^2) \Big|_{t=0} = \text{tr}(N^2)$$

Continuing,  $\text{tr}(N^k) = 0$  for any  $k \in \mathbb{N}$ .

(we need only  $k=1, \dots, n$ )

Now, suppose for contradiction, that  $N$  is not nilpotent. (6)

Then  $N$  has some non-zero eigenvalues  $\lambda_1, \dots, \lambda_r$ . \* If all eig = 0 then nilpotent

Let  $m_i$  be the multiplicity of  $\lambda_i \in \mathbb{C}$ . Then

$$\begin{cases} m_1 \lambda_1 + \dots + m_r \lambda_r = 0 \\ \vdots \\ m_1 \lambda_1^r + \dots + m_r \lambda_r^r = 0 \end{cases}$$

\* Schur decomposition:

$A \in \mathbb{C}^{n \times n}$  then  $A = UTU^{-1}$   
 $T$  upper triangular, diagonal are eigenvalues

Suppose that  $\lambda = 0$ . Then  $T^n = 0$ . Thus  $A^n = UT^nU^{-1} = 0$

Thus

(\*)

$$\begin{pmatrix} \lambda_1 & \lambda_2 & \lambda_3 & \dots & \lambda_r \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 & \dots & \lambda_r^2 \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ \lambda_1^r & \lambda_2^r & \lambda_3^r & \dots & \lambda_r^r \end{pmatrix} \begin{pmatrix} m_1 \\ \vdots \\ \vdots \\ m_r \end{pmatrix} = \begin{pmatrix} 0 \\ \vdots \\ \vdots \\ 0 \end{pmatrix}$$

But this is a Vandermonde matrix, so  $\prod_{0 \leq i < j \leq r} (\lambda_i - \lambda_j)$

$$\det \begin{pmatrix} \lambda_1 & \lambda_2 & \lambda_3 & \dots & \lambda_r \\ \lambda_1^2 & \lambda_2^2 & \lambda_3^2 & \dots & \lambda_r^2 \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ \lambda_1^r & \lambda_2^r & \lambda_3^r & \dots & \lambda_r^r \end{pmatrix} = \lambda_1 \dots \lambda_r \det \begin{pmatrix} 1 & 1 & \dots & 1 \\ \lambda_1 & \lambda_2 & \dots & \lambda_r \\ \vdots & \vdots & \ddots & \vdots \\ \lambda_1^{r-1} & \dots & \dots & \lambda_r^{r-1} \end{pmatrix}$$

$\neq 0$  since  $\lambda_i \neq \lambda_j$  for  $i \neq j$

Thus (\*) has unique solution:  $m_i = 0$ . Contradiction.

For the other direction, suppose that  $N$  is nilpotent. Then  $N^k = 0$  for some  $k \in \mathbb{N}$ . ⑦

We claim

$$\det(I + N) = 1$$

Suppose, for contradiction, that  $\det(I + N) \neq 1$ .

Then  $I + N$  must have a non-unity eigenvalue

$$\lambda u = (I + N)u = u + Nu \quad u \neq 0.$$

Thus  $Nu = (\lambda - 1)u$

Note that  $(\lambda - 1)u \neq 0$ . Thus, by induction:

$$N^k u = (\lambda - 1)^k u \quad \text{for all } k \in \mathbb{N}.$$

This contradicts  $N$  being nilpotent. Thus all  $\lambda = 1$ .

Note, since  $I + N$  has all eigenvalues one bigger than  $N$ , it follows that all eigenvalues of  $N$  are zero.

# Return to Asymptotic geodesics

⑧

Recall  $X_t(a) = a + t u_0(a)$

$$\nabla X_t(a) = I + t A_0 \quad (A_0 = \nabla_a u_0(a))$$

we must have

$$\det(I + t A_0) = 1 \quad \text{for all } t \in \mathbb{R}$$

This holds iff  $A_0$  is Nilpotent. In

which case

$$(I + t A_0)^{-1} = \sum_{m=0}^k (-A_0)^m \quad (A_0^k = 0)$$

$$= I - A_0 + A_0^2 - A_0^3 + \dots + (-A_0)^k$$

Using,

$$\nabla_x u(a + t u_0(a), t) = (I + t \nabla_a u_0(a))^{-1} \nabla_a u_0(a)$$

We see

$$(\nabla_x u)(X_t(a), t) = \sum_{m=0}^{d-2} (-t)^m (\nabla_a u_0)^{m+1}$$

In dim  $d$ , grows  
like  
 $|\nabla u|_{\infty} \sim t^{d-2}$



# Examples of Nilpotent data

(19)

$$d=2 \quad u_0(a_1, a_2) = \begin{pmatrix} v(a_2) \\ 0 \end{pmatrix}$$

$$\nabla_a u_0(a) = \begin{pmatrix} 0 & v'(a_2) \\ 0 & 0 \end{pmatrix}$$

$$d=3 \quad u_0(a_1, a_2, a_3) = \begin{pmatrix} u_1(a_2, a_3) \\ u_2(a_3) \\ 0 \end{pmatrix}$$

$$\nabla u_0(a) = \begin{pmatrix} 0 & \partial_2 u_1 & \partial_3 u_1 \\ 0 & 0 & \partial_3 u_2 \\ 0 & 0 & 0 \end{pmatrix}$$

In this case

$$\dot{x}_1(t) = u_1(x_2, x_3)$$

$$\dot{x}_2(t) = u_2(x_3)$$

$$\dot{x}_3(t) = 0$$

$$x_1(t) = a_1 + u_1(a_2 + t u_2(a_3), a_3)$$

$$\Rightarrow x_2(t) = a_2 + t u_2(a_3)$$

$$x_3(t) = a_3$$

$$\frac{d}{dt} u_1(a) = \dot{x}_2 \partial_2 u_1 + \dot{x}_3 \partial_3 u_1 = u_2(x_3) \cdot \nabla_2 u_1(x_2, x_3)$$

$\partial_t u_1'(x_2, x_3)$

Questions:

① In dimension  $d=2$ , classify all pressureless solutions. Are they all steady? If steady, are they all shear flows? At what regularity?

② One can construct solutions that blow-up in finite time, in infinite space dim.

Take data

$$u_0(x) = \begin{pmatrix} u_1(a_2) \\ u_2(a_3) \\ u_3(a_4) \\ \vdots \\ u_{d-1}(a_d) \\ 0 \end{pmatrix}$$

Note:  $\|\nabla u_0\|_{L^\infty} := \sup_x \max_{1 \leq i \leq d} \sum_{j=1}^d |(\nabla u_0)_{ij}| < \infty$  ↖ ind of d.

On other hand

$$\|\nabla u(0, t)\|_{L^\infty} = \sum_{i=0}^d t^i = \frac{1-t^{d+1}}{1-t}$$

as  $d \rightarrow \infty$ , this blows up at time  $\underline{1!}$

# Singularity formation

(11)

Let's focus on  $d=2$  and  $M=\mathbb{R}$ . We have

$$u_x(x_f(a), t) = \frac{u_0'(a)}{1 + u_0'(a)}$$

Thus the first singularity emerges from the label  $a_*$  at which  $u_0'(a_*)$  is most negative.

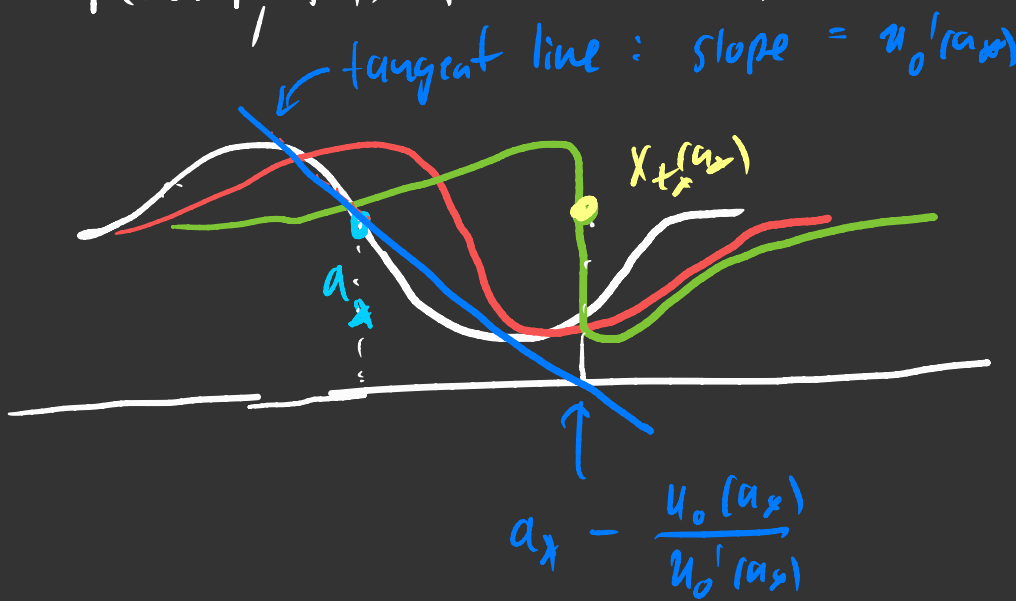
The time of blowup is explicit

$$t_* = -\frac{1}{u_0'(a_*)}$$

the location is also

$$\begin{aligned} X_{t_*}(a_*) &= a_* + t_* u_0(a_*) \\ &= a_* - \frac{u_0(a_*)}{u_0'(a_*)} \end{aligned}$$

Remark; this formula looks like Newton's method!



Note that for  $|a - a_*| \ll 1$ , we have

(12)

$$\begin{aligned}
 u_0'(a) &= u_0'(a_*) && (< 0, \text{ by assumption}) \\
 &+ \cancel{u_0''(a_*)} (a - a_*) && \text{since } a_* \text{ is minimum of } u_0'(a_*) \\
 &+ C (a - a_*)^2 && (C > 0, \text{ since minimum}) \\
 &+ \mathcal{O}((a - a_*)^3) && (\text{assuming } u_0 \in C^4)
 \end{aligned}$$

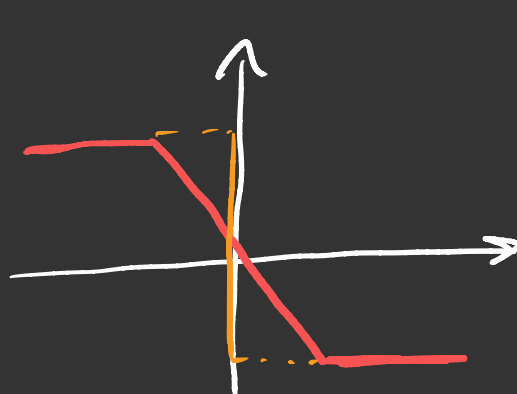
Thus, at time  $t_*$ ,

$u_x(x_{t_*}(a), t_*)$  is finite  $\forall a \neq a_*$ .

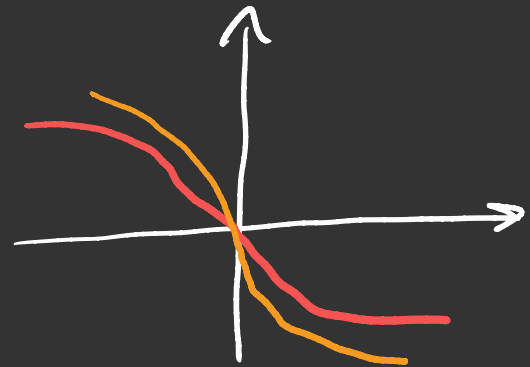
First singularity occurs at a single point in spacetime

Define: A data is a global minimum and generic if

- $u_0 \in C^4$
- $u_0'(a_*) < 0$
- $u_0''(a_*) = 0$
- $u_0'''(a_*) > 0$



non-generic: jump



generic: cusp

What regularity should we expect?

(13)

Consider data

$$u_0(a) = -a + a^n, \quad n \geq 3, \quad n \text{ odd}$$

Note that  $a_* = 0$  is global minimum

$$u_0'(0) = -1 < 0$$

$$u_0''(0) = 0$$

$$u_0'''(0) = n \cdot (n-1) \cdot (n-2) a^{n-3}$$

$$= \begin{cases} 6 & \text{if } n=3 \\ 0 & \text{if } n \geq 3 \end{cases}$$

generic  
non-generic

Time of blowup is

$$t_* = -\frac{1}{u_0'(a_*)} = 1$$

The flowmap is then is

$$X_{t_*}(a) = a + u_0(a) = a^n$$

Recalling that

$$u(X_t(a), t) = u_0(a)$$

at time  $t_n = 1$ , we have

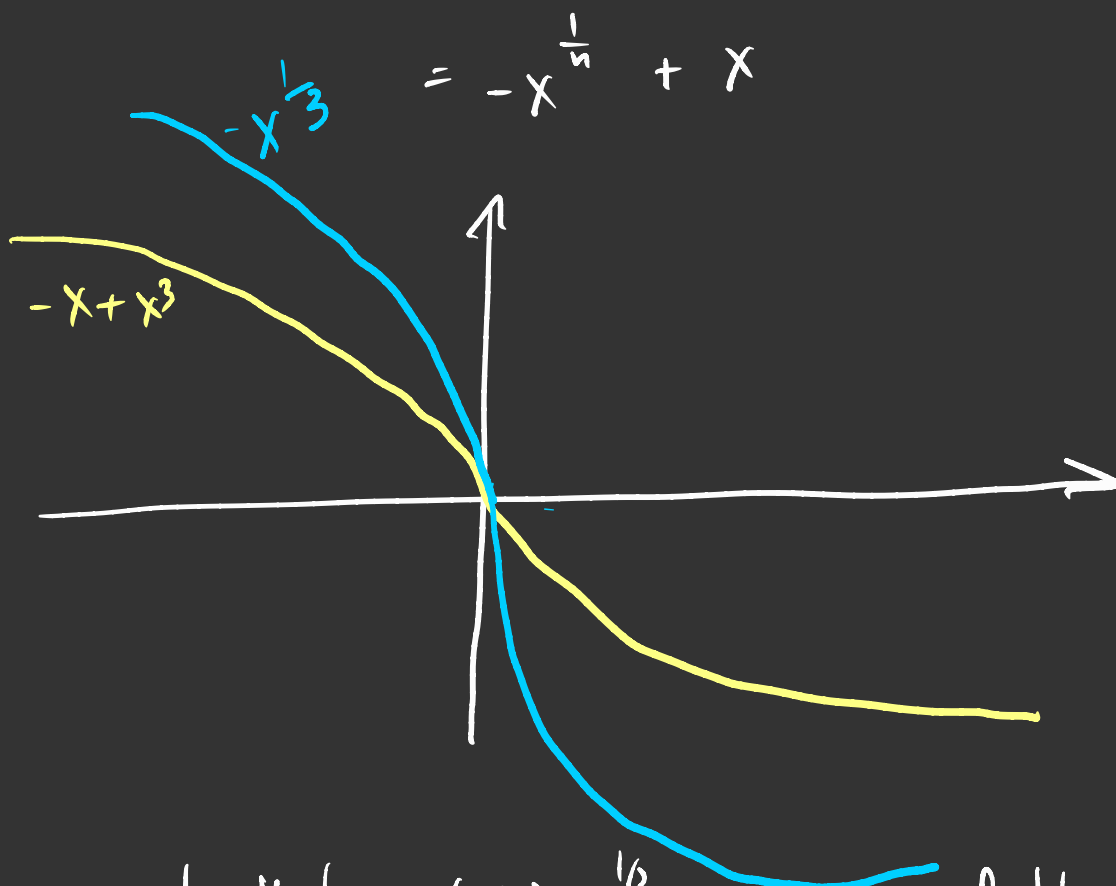
$$u(a^n, 1) = u_0(a)$$

Inverting the homeomorphism  $a^n$  gives

$$X_{t_n}^{-1}(x) = x^{\frac{1}{n}}.$$

Thus

$$u(x, 1) = u_0(x^{\frac{1}{n}})$$



Thus, we expect that  $u(\cdot, t_n) \in C^{\frac{1}{3}}$  at time of blowup!

# Self-similarity

Note that if

$$X_t(a) = a + t u_0(a)$$

then  $y = A_t(x)$ , the inverse, solves

$$x = y + t u_0(y).$$

For the special data:  $x = (1-t)y + t y^3$

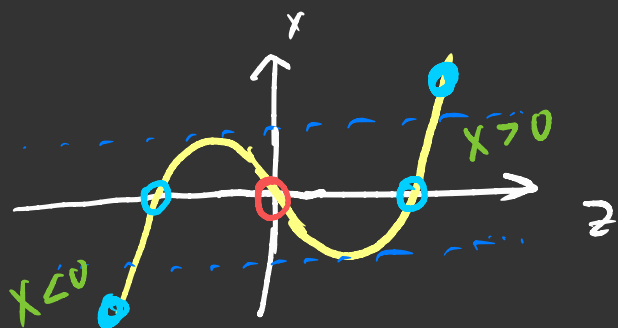
Substitute  $z = \left(\frac{t}{1-t}\right)^{1/2} y$ . Then

$$t y^3 = \frac{(1-t)^{3/2} z^3}{\sqrt{t}} \quad (1-t)y = \frac{(1-t)^{3/2} z}{\sqrt{t}}$$

$$z^3 + z = \frac{\sqrt{t}}{(1-t)^{3/2}} x \equiv X$$

Thus

$$z = F(X), \quad \begin{array}{l} F \text{ explicit} \\ F'(0) = -1 \end{array}$$



In fact  $z^3 + z \equiv X$  is solved by (16)

$$z = F(x) = \left[ \frac{x}{2} + \left( \frac{1}{27} + \frac{x^2}{4} \right)^{1/2} \right]^{1/3} - \left[ -\frac{x}{2} + \left( \frac{1}{27} + \frac{x^2}{4} \right)^{1/2} \right]^{1/3}$$

Thus

$$\begin{aligned} u(x,t) &= u_0(A_t(x)) \\ &= -A_t(x) + A_t(x)^3 \\ &= -\frac{(1-t)^{1/2}}{\sqrt{t}} F\left(\frac{\sqrt{t} x}{(1-t)^{3/2}}\right) + \frac{(1-t)^{3/2}}{t^{3/2}} F^3\left(\frac{\sqrt{t} x}{(1-t)^{3/2}}\right) \end{aligned}$$

and

$$u_x(x,t) = \underbrace{\frac{1}{1-t} F'\left(\frac{\sqrt{t} x}{(1-t)^{3/2}}\right)}_{\text{blowing up at } t=1} + \underbrace{3(F'F)\left(\frac{\sqrt{t} x}{(1-t)^{3/2}}\right)}_{\text{bounded as } t \rightarrow 1}$$

blowing up at  
 $t=1$   
in a self-similar  
way