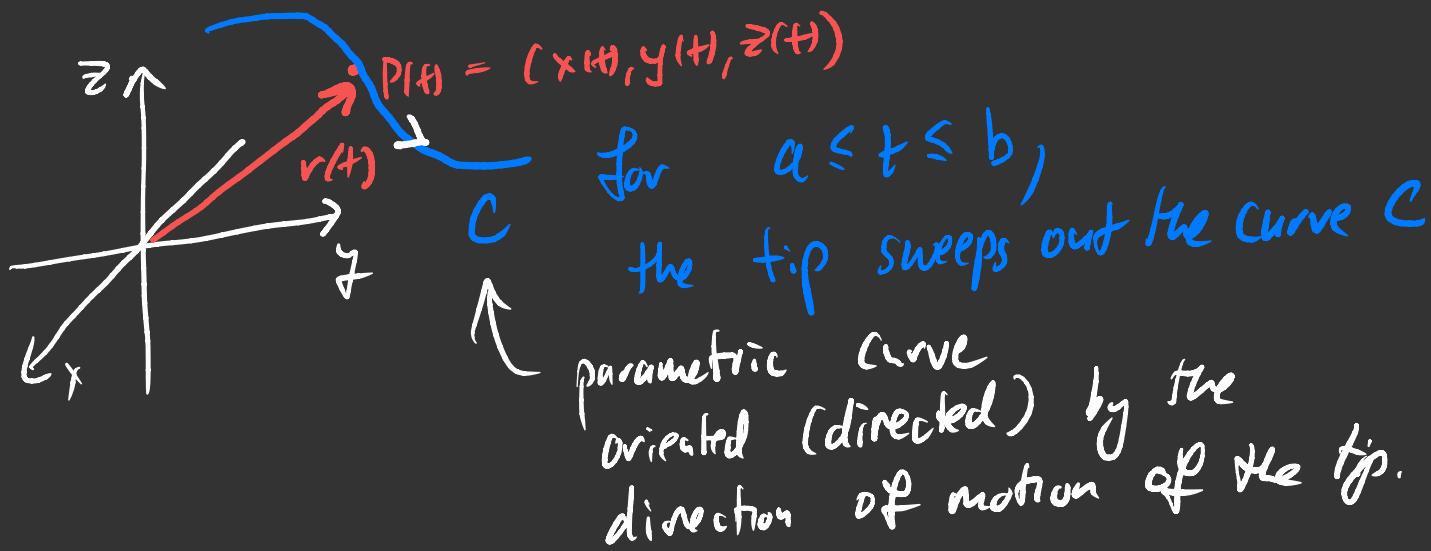


Vector functions : function whose values are vectors

$$\vec{r}(t) = (x(t), y(t), z(t))$$



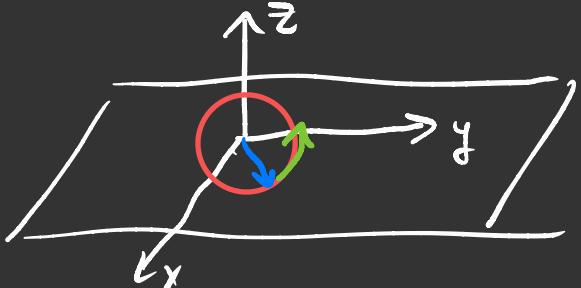
Example: $\vec{r}(t) = \vec{r}_0 + t\vec{u}$. : motion along straight line with constant speed.

Example: $\vec{r}(t) = (\cos t, \sin t, 0)$

- planar motion

- note that $(x(t))^2 + (y(t))^2 = (\cos t)^2 + (\sin t)^2 = 1$

thus this motion remains on unit circle



motion is counter clockwise observed from tip of \vec{R} .

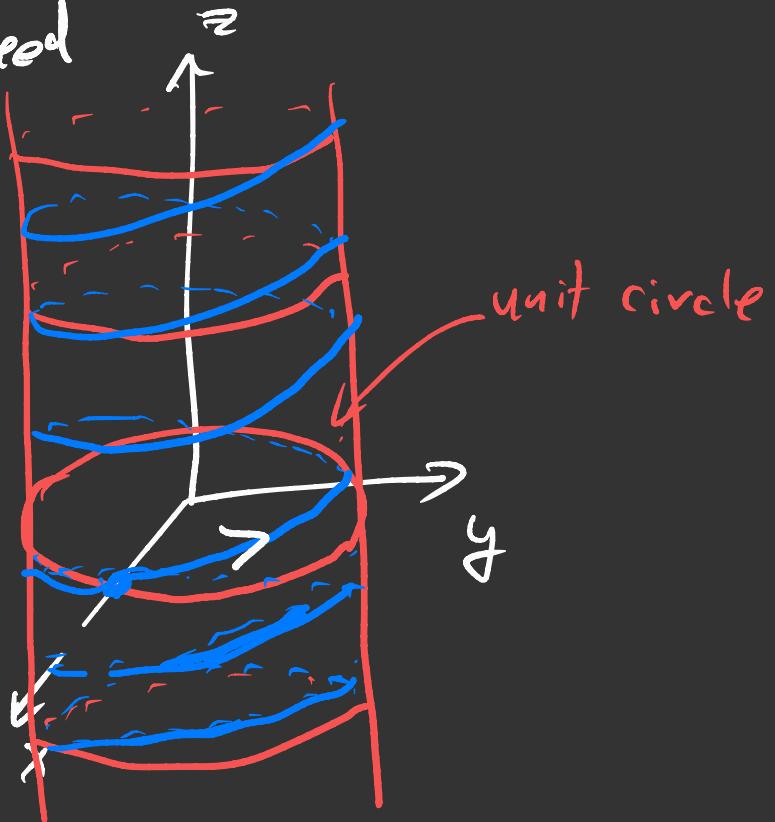
Example : $\mathbf{r}(t) = (\cos(t), \sin(t), t)$

• $(x(t))^2 + (y(t))^2 = 1$

Moving up in z at const speed

Helix

ccw motion
with respect to \vec{k} .



Differentiation

$$\vec{r}(t) = (x(t), y(t), z(t))$$

$$\vec{r}'(t) = (x'(t), y'(t), z'(t))$$

$$\vec{r}'(t) = \lim_{h \rightarrow 0} \frac{\vec{r}(t+h) - \vec{r}(t)}{h}$$

Properties of derivative

1) Linearity:

$$(\mathbf{r}_1(t) + \mathbf{r}_2(t))' = \mathbf{r}_1'(t) + \mathbf{r}_2'(t)$$

2) $k \in \mathbb{R}$, $(k\mathbf{r}(t))' = k\mathbf{r}'(t)$.

Proof: Limit definition.

3) Leibnitz Rules

Scalar a) $(f(t)g(t))' = f'(t)g(t) + f(t)g'(t)$

Scalar multiple b) $(k(t)\vec{r}(t))' = k'(t)\vec{r}(t) + k(t)\vec{r}'(t)$

dot product c) $(\vec{r}_1(t) \cdot \vec{r}_2(t))' = \vec{r}_1'(t) \cdot \vec{r}_2(t) + \vec{r}_1(t) \cdot \vec{r}_2'(t)$

cross product d) $(\vec{r}_1(t) \times \vec{r}_2(t))' = \vec{r}_1'(t) \times \vec{r}_2(t) + \vec{r}_1(t) \times \vec{r}_2'(t)$

all are proved in some way. Let's prove (d).

$$\begin{aligned} \text{Write } (\vec{r}_1(t) \times \vec{r}_2(t))' &= \lim_{h \rightarrow 0} \frac{\vec{r}_1(t+h) \times \vec{r}_2(t+h) - \vec{r}_1(t) \times \vec{r}_2(t)}{h} \\ &= \lim_{h \rightarrow 0} \frac{(\vec{r}_1(t+h) - \vec{r}_1(t)) \times \vec{r}_2(t+h) - \vec{r}_1(t) \times (\vec{r}_2(t+h) - \vec{r}_2(t))}{h} \\ &= \lim_{h \rightarrow 0} \left[\left(\frac{\vec{r}_1(t+h) - \vec{r}_1(t)}{h} \right) \times \vec{r}_2(t+h) + \vec{r}_1(t) \times \left(\frac{\vec{r}_2(t+h) - \vec{r}_2(t)}{h} \right) \right] \end{aligned} \quad (3)$$

3)

Chain Rule

$$\vec{r}(t) \quad \text{and} \quad t = t(s)$$

$$(\vec{r}(t(s)))' = t'(s) \vec{r}'(t(s))$$

Example: $r(t) = (\cos(t), \sin(t), t)$

$$t(s) = s^2$$

$$r(t(s)) = (\cos(s^2), \sin(s^2), s^2)$$

$$\begin{aligned} \frac{d}{ds} r(t(s)) &= \frac{d}{ds} (\cos(s^2), \sin(s^2), s^2) \\ &= (-\sin(s^2) \cdot (2s), \cos(s^2) \cdot (2s), 1) \end{aligned}$$

$$= 2s (-\sin(s^2), \cos(s^2), 1)$$

$$= \frac{dt}{ds} \cdot \frac{d\vec{r}}{dt}(t(s))$$

$$\frac{dt}{ds} = 2s.$$

✓

4) Higher derivatives

$$\vec{r}(t) = (x(t), y(t), z(t))$$

$$\vec{r}'(t) = (x'(t), y'(t), z'(t))$$

$$\vec{r}''(t) = (x''(t), y''(t), z''(t))$$

⋮

$$\vec{r}^{(n)}(t) = (x^{(n)}(t), y^{(n)}(t), z^{(n)}(t))$$

Example: $(\vec{r}_1(t) \cdot \vec{r}_2(t))^I = \vec{r}_1'(t) \cdot \vec{r}_2(t) + \vec{r}_1(t) \cdot \vec{r}_2'(t)$

$$(\vec{r}_1(t) \cdot \vec{r}_2(t))^H = \vec{r}_1''(t) \cdot \vec{r}_2(t) + 2\vec{r}_1'(t) \cdot \vec{r}_2'(t) + \vec{r}_1(t) \cdot \vec{r}_2''(t)$$

Rate of change of volume of a parallelepiped.

Consider

$$\vec{a}(t), \vec{b}(t) \text{ and } \vec{c}(t).$$

$$V(\vec{a}(t), \vec{b}(t), \vec{c}(t)) = (\vec{a}(t) \times \vec{b}(t)) \cdot \vec{c}(t)$$

$$\begin{aligned} \frac{d}{dt} V(\vec{a}(t), \vec{b}(t), \vec{c}(t)) &= \frac{d}{dt} \left((\vec{a}(t) \times \vec{b}(t)) \cdot \vec{c}(t) \right) \\ &= (\vec{a}'(t) \times \vec{b}(t)) \cdot \vec{c}(t) \\ &\quad + (\vec{a}(t) \times \vec{b}'(t)) \cdot \vec{c}(t) \\ &\quad + (\vec{a}(t) \times \vec{b}(t)) \cdot \vec{c}'(t) \\ &= V(\vec{a}'(t), \vec{b}(t), \vec{c}(t)) \\ &\quad + V(\vec{a}(t), \vec{b}'(t), \vec{c}(t)) \\ &\quad + V(\vec{a}(t), \vec{b}(t), \vec{c}'(t)) \end{aligned}$$

Integral of vector function

Primitive :

$$\vec{r}(t) = (x(t), y(t), z(t))$$

Def: $\vec{R}(t)$ is called the primitive if $\vec{R}'(t) = \vec{r}(t)$.

$$\vec{R}(t) = (\vec{u}(t), \vec{v}(t), \vec{w}(t))$$

$$\vec{R}'(t) = (\vec{u}'(t), \vec{v}'(t), \vec{w}'(t))$$

Then

$$\vec{u}'(t) = \vec{x}(t)$$

$$\vec{v}'(t) = \vec{y}(t)$$

$$\vec{w}'(t) = \vec{z}(t)$$

Then

$$\vec{u}(t) = \int \vec{x}(t) dt + C_1$$

$$\vec{v}(t) = \int \vec{y}(t) dt + C_2$$

$$\vec{w}(t) = \int \vec{z}(t) dt + C_2$$

or

$$\vec{R}(t) = \int \vec{r}(t) dt + \vec{C}$$

$$\underline{\text{Ex:}} \quad \vec{r}(t) = (\cos t, \sin t, t)$$

$$\begin{aligned} \vec{R}(t) &= \left(\int \cos t dt + C_1, \int \sin t dt + C_2, \int t dt + C_3 \right) \\ &= (\sin t, -\cos t, \frac{t^2}{2}) + (C_1, C_2, C_3) \end{aligned}$$

Definite integral:

$$\begin{aligned} \int_a^b \vec{r}(t) dt &= \vec{R}(b) - \vec{R}(a) \\ &= \left(\int_a^b x(t) dt, \int_a^b y(t) dt, \int_a^b z(t) dt \right) \end{aligned}$$

$$\begin{aligned} \underline{\text{Ex:}} \quad &\int_0^{\pi/4} (\cos t, \sin t, t) dt \\ &= \left(\frac{1}{\sqrt{2}}, 1 - \frac{1}{\sqrt{2}}, \frac{1}{2} \left(\frac{\pi}{4}\right)^2 \right) \end{aligned}$$

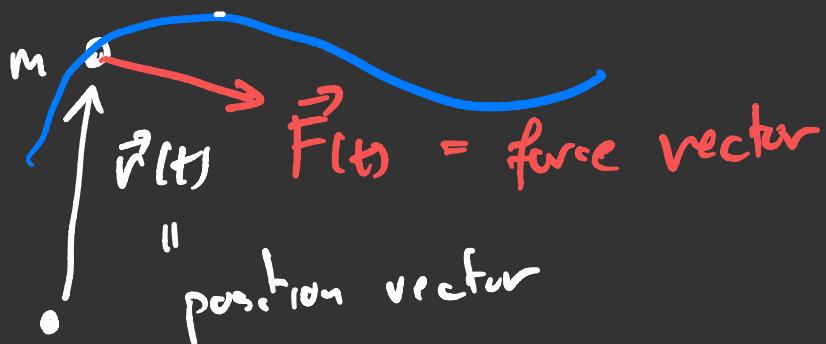
Since

$$\int_0^{\pi/4} \cos t dt = \sin\left(\frac{\pi}{4}\right) - \sin(0) = \frac{1}{\sqrt{2}}$$

$$\int_0^{\pi/4} \sin t dt = -\cos\left(\frac{\pi}{4}\right) + \cos(0) = 1 - \frac{1}{\sqrt{2}}$$

Some dynamical problems.

① 2d Newton's Law : mass m



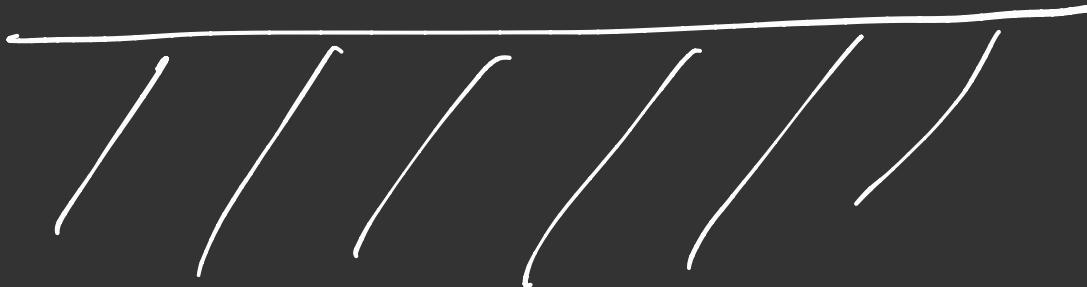
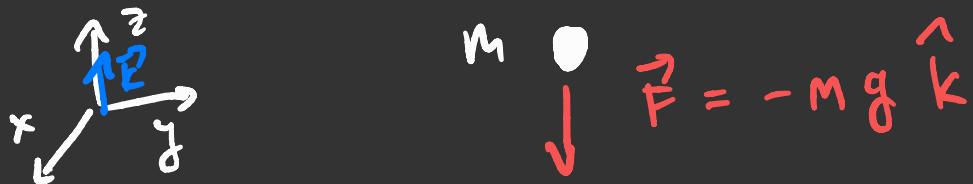
$$\vec{F}(t) = m \vec{a}(t) = m \vec{r}''(t)$$

↑
acceleration.

Vectorial formulation of Newton's Law

Examples : Motion in gravitational field

Assume the earth is horizontal and flat, and the force of gravity acts vertically:



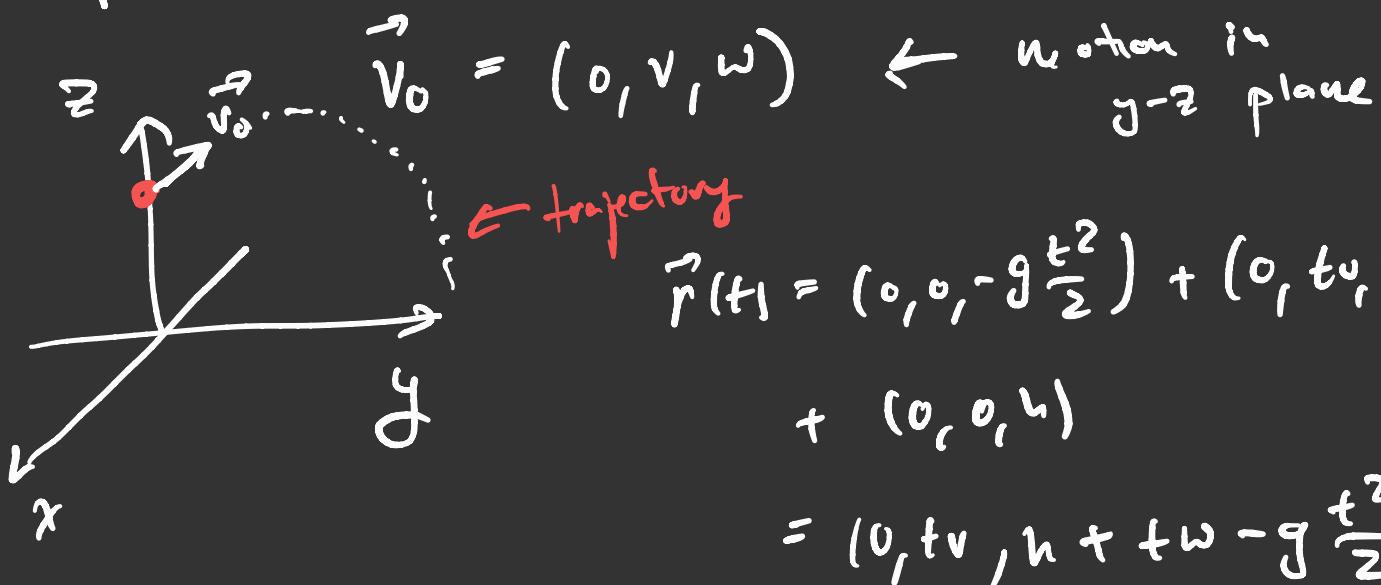
$$\vec{a} = \frac{\vec{F}}{m} = -g \hat{k} = (0, 0, -g)$$

$$\begin{aligned}\vec{a} = \frac{d}{dt} \vec{v}(t) &\Rightarrow \vec{v}(t) = \int \vec{a}(t) dt + \vec{v}_0 \\ &= - \int g \hat{k} dt + \vec{v}_0 \\ &= -gt \hat{k} + \vec{v}_0\end{aligned}$$

$$\begin{aligned}\vec{v}(t) &= \frac{d}{dt} \vec{r}(t) \\ \vec{r}(t) &= \int \vec{v}(t) dt + \vec{r}_0 \\ &= -g \frac{t^2}{2} \hat{k} + t \vec{v}_0 + \vec{r}_0\end{aligned}$$

$$\boxed{\vec{r}(t) = \vec{v}_0 + t \vec{v}_0 - g \frac{t^2}{2} \hat{k}}$$

Example: $\vec{r}_0 = (0, 0, H)$ height = H



Let's find equation for the trajectory in the y-z plane, e.g. $z = z(y)$

Note $y = t v \Rightarrow t = \frac{y}{v}$

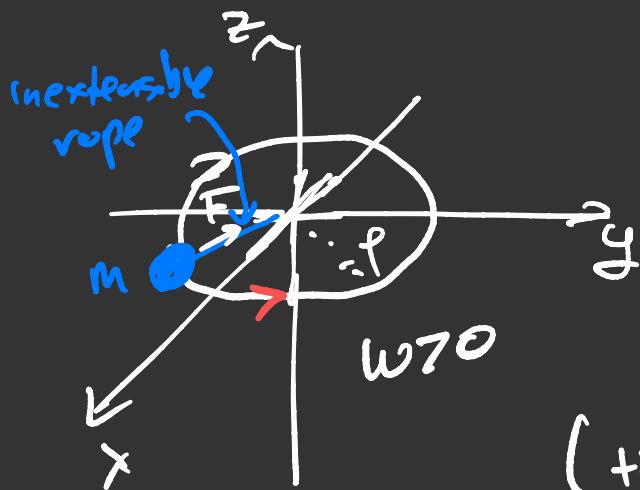
$$z = H + t w - g \frac{t^2}{2}$$

$$= H + \frac{w}{v} y - \frac{g}{2v^2} y^2$$

upside-down parabola.

2) Circular motion

$r > 0 \leftarrow$ radius
 $\omega = \text{frequency}$



$$\vec{r}(t) = (\rho \cos \omega t, \rho \sin \omega t, 0)$$

$$\text{period } T = \frac{2\pi}{\omega}$$

(time it takes to make one turn)

Rope pulls body to the center, so
 force \vec{F} is directed in. Let's find \vec{F} :

$$\vec{F} = m \vec{a} = m \vec{r}''(t)$$

while

$$\vec{r}'(t) = \omega (-\rho \sin(\omega t), \rho \cos(\omega t), 0)$$

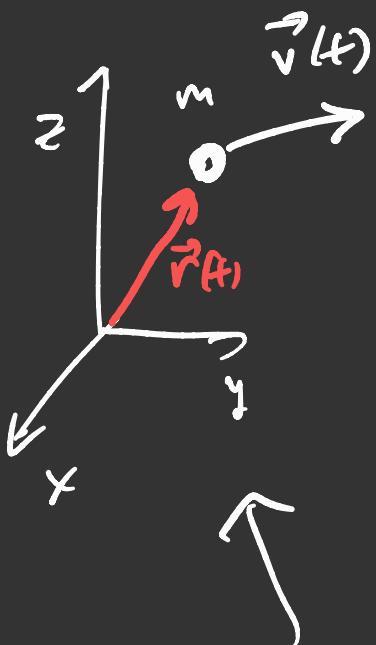
$$\begin{aligned} \vec{r}''(t) &= -\omega^2 (\rho \sin(\omega t), \rho \cos(\omega t), 0) \\ &= -\omega^2 \vec{r}(t). \end{aligned}$$

Thus

$$\vec{F}(t) = -m \omega^2 \vec{r}(t)$$

centrifugal force

3) Angular momentum and torque.



Angular Momentum

$$\vec{M} = m \vec{r} \times \vec{v}$$

perpendicular to \vec{r}, \vec{v}
and directed so that
 $(\vec{r}, \vec{v}, \vec{M})$ form a right triple

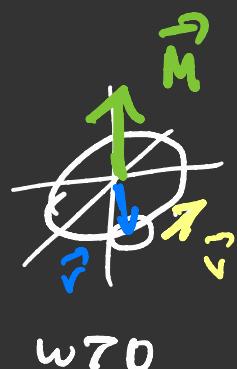
In this case, \vec{M} points inside the page.

In previous example

$$\vec{r} = (\rho \cos(\omega t), \rho \sin(\omega t), 0)$$

$$\vec{v} = (-\rho \sin(\omega t), \rho \cos(\omega t), 0)$$

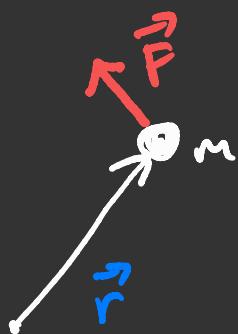
$$\vec{M} = m \vec{r} \times \vec{v} = m \begin{vmatrix} \vec{i} & \vec{j} & \vec{k} \\ \rho \cos \omega t & \rho \sin \omega t & 0 \\ -\rho \sin \omega t & \rho \cos \omega t & 0 \end{vmatrix}$$



$$= m \begin{vmatrix} \rho \cos(\omega t) & \rho \sin(\omega t) & \vec{k} \\ -\rho \sin(\omega t) & \rho \cos(\omega t) & \end{vmatrix}$$

$$= m \rho^2 \omega (\cos^2(\omega t) + \sin^2(\omega t)) \vec{k} = m \rho^2 \omega \vec{k}$$

Torque



Torque

$$\vec{\tau} = \vec{r} \times \vec{F}$$

If a force is acting on a body, it changes its angular momentum.
How does it change? Analogue of 2nd Newton's Law.

Theorem : If the mass is moving along a trajectory $\vec{r}(t)$ under the influence of force \vec{F} , then

$$\frac{d}{dt} \vec{M}(t) = \vec{\tau}(t)$$

where $\vec{M} = m \vec{r} \times \vec{v}$ is the angular momentum.

(Mass) \times Velocity = linear momentum.

\vec{M} is angular momentum

Proof:

$$\vec{M} = \vec{r} \times (m \vec{v})$$

$$\begin{aligned}
 \frac{d}{dt} \vec{M}(t) &= m \vec{r}' \times \vec{v} + m \vec{r} \times \vec{v}' \\
 &= m \vec{v} \times \vec{v} \quad \text{(cancel)} + m \vec{r} \times \vec{r}'' \\
 &= \vec{r} \times (m \vec{r}'') \\
 &= \vec{r} \times \vec{F} \\
 &= \vec{\tau}.
 \end{aligned}$$

(Leibnitz)
 $\vec{a} \times \vec{a} = 0$
 for all \vec{a})

Pf: change places
 $\vec{u} \times \vec{w} = -\vec{w} \times \vec{u}$

□

Remarkable particular case:

If $\vec{\tau}(t) = 0$, then $\vec{M}(t) = \text{const.}$

($\frac{d}{dt} \vec{M}(t) = \vec{0}$ so $\vec{M}(t) = \vec{M}_0$)

What's an example of this? Central force

$$\vec{F}(t) = k(t) \vec{r}(t)$$

$$\begin{aligned}
 \vec{\tau} &= \vec{r} \times \vec{F} \\
 &= k \vec{r} \times \vec{r} = 0.
 \end{aligned}$$

Example:

Inverse square law

$$\|\vec{F}\| \sim \|\vec{r}\|^{-2}$$

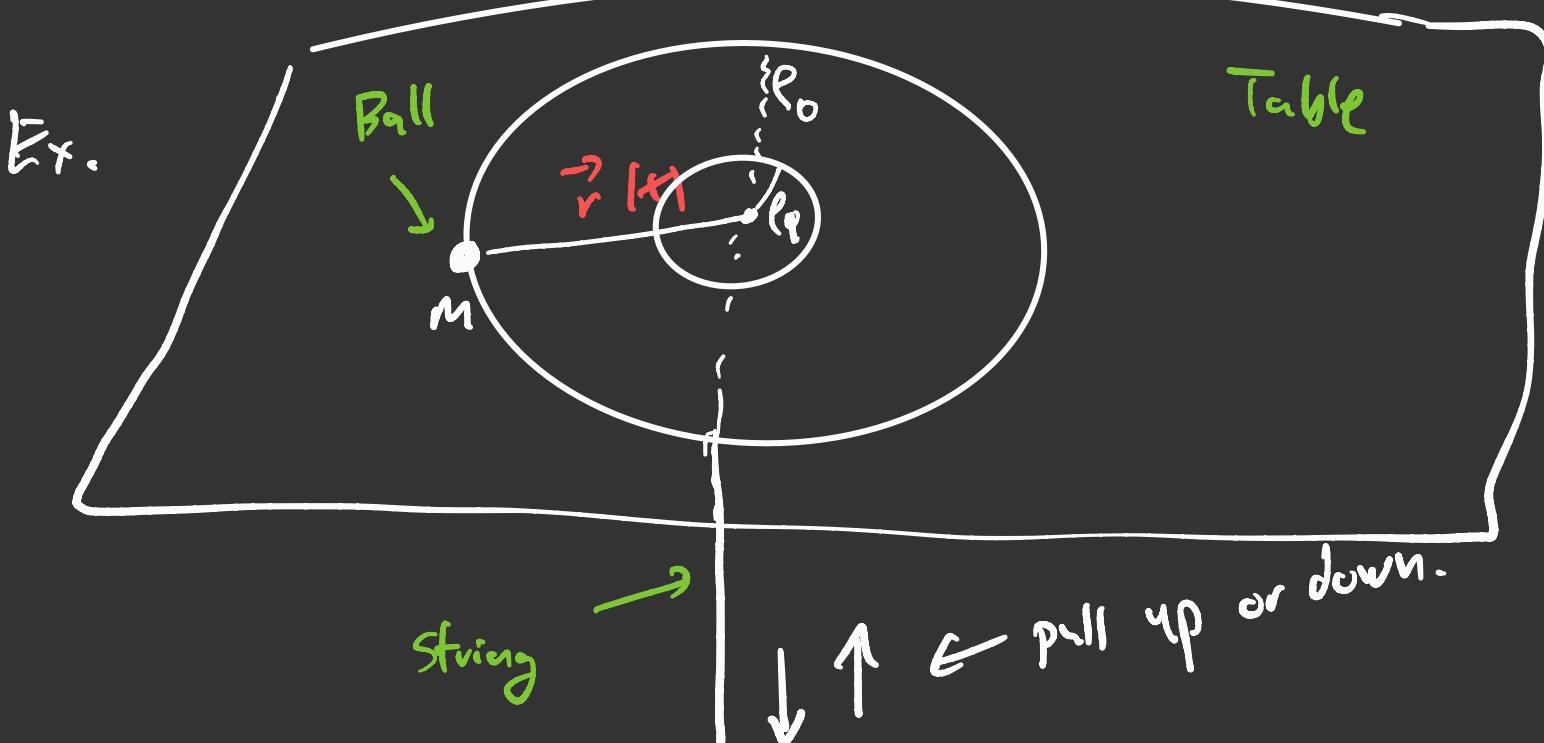
Plane & M_p

$\vec{F} = -\frac{GM_s M_p}{\|\vec{r}\|^3} \vec{r}$

Sun, M_s

⑦

Thus, if \vec{F} is a central force, then
 $\vec{M} = m \vec{r} \times \vec{v}$ is constant.



Thus we can prescribe $\|\vec{r}(t)\|$ arbitrarily.

$$\vec{M} = m \vec{r}(t) \times \vec{v}(t) = \text{constant}$$

$$\vec{M} = m r^2 \vec{\omega} \hat{k} \quad \text{angular momentum.}$$

$$\vec{M}_0 = m r_0^2 \vec{\omega}_0 \hat{k} \quad \vec{M}_1 = m r_1^2 \vec{\omega}_1 \hat{k}$$

$$\Rightarrow \frac{\omega_1}{\omega_0} = \frac{r_0^2}{r_1^2} \Leftrightarrow$$

$$\boxed{\omega_1 = \frac{r_0^2}{r_1^2} \omega_0}$$

If you reduce length by half, $r_1 = \frac{1}{2} r_0$.
then $\omega_1 = 4 \omega_0$ and $\|v_1\| = 2 \|v_0\|$.